## 35, 1957 <br> 

; a function of temperature at $\geq$ been normalized to agree at serlapping.
esistance of rubidium at ber 3 ) show this anomaly e evident, but in sample 4 1 deviation from linearity $\mathrm{t} 150^{\circ} \mathrm{K}$. Samples 1 and 2 hood of the anomaly (cf. alling temperature being implify the diagram only
shows little sign of the since at this temperature the resistance increases lgcock 1956) confirm that be suppressed.
and Pearson (1955) have sorigin is still obscure.

PRESSURE EFFECTS
The isothermal change in resistance with pressure was measured in the neighborhood of room temperature on samples 1 and 2. As might be expected with such a fine capillary the change in resistance of sample 1 for a given pressure was appreciably lower than that in sample 2. Even with sample 2 the resistivity change at 2500 atm . was about $10 \%$ smaller than that deduced from Bridgman's measurements (Bridgman 1946) at higher pressures. (Bridgman's measurements refer to relative resistance changes of a given sample and have to be transformed into resistivity changes using the compressibility.) This difference is presumably due to the fact that we are using specimens melted into capillaries. The present experiments serve nevertheless to demonstrate that the general method is convenient and to give at least a qualitative picture of the pressure dependence of resistance down to very low temperatures.

In describing these results, it is convenient to divide the temperature range as follows:
(1) measurements between $50^{\circ} \mathrm{K}$. and room temperature,
(2) measurements between $4^{\circ} \mathrm{K}$. and $50^{\circ} \mathrm{K}$.,
(3) measurement of the residual resistance.


Fig. 3. The resistivity of a rubidium sample (sample 2) at two pressures in the higher temperature range.

